

Application of neural networks for computing path integrals in quantum theory

Artyom Vasiliev¹, Alexandr Ivanov¹,
Dmitry Salnikov^{1,2} and Vsevolod Chistiakov¹

salnikov@inr.ru

¹Lomonosov Moscow State University

²Institute for Nuclear Research of the Russian Academy of Science

The work of DS and VC is supported by «BASIS» Foundation grants:

№ 23-2-2-40-1 and № 23-2-9-18-1.

June 19, 2024

Introduction

- Path integral in imaginary time:

$$\langle O \rangle_\beta = \frac{1}{Z} \int D[x(t)] O[x(t)] e^{-S_E[x(t)]}, \quad x(0) = x(\beta). \quad (1)$$

- Gibbs average:

$$\langle O \rangle_\beta = \frac{1}{Z} \text{Tr} [e^{-\beta H} O] = \frac{1}{Z} \sum_n \langle n | O | n \rangle e^{-\beta E_n}. \quad (2)$$

- Low temperature limit: $\lim_{\beta \rightarrow \infty} \langle O \rangle_\beta = \langle 0 | O | 0 \rangle$.

- Theory on the lattice:

$$D[x(t)] = \prod_n dx_n \equiv d^n x, \quad S_E[x(t)] = S_E(x_1, \dots, x_n) \equiv S_E(x). \quad (3)$$

- QFT generalization: $t \rightarrow x^\mu, x(t) \rightarrow \phi(x)$.

Monte Carlo calculation

- Average over a distribution:

$$\langle O \rangle_{\beta} = \int d^n x O(x) p(x) \simeq \frac{1}{N} \sum_{j=1}^N O(x^{(j)}), \quad (4)$$

where

$$\{x^{(j)}\} \sim p(x) = \frac{1}{Z} e^{-S_E(x)}. \quad (5)$$

- Path integral calculation = generation of the samples $\{x^{(j)}\}$ with target density distribution function.

Markov chain Monte Carlo (Metropolis)

- Starting with a sample of the «cold» trajectories:

$$X_0 = \{x^{(j)}\}, \quad x_i^{(j)} = 0. \quad (6)$$

- Random modification of the trajectories:

$$\tilde{x}_i^{(j)} = x_i^{(j)} + \sigma(\tau)u_i^{(j)}, \quad u_i^{(j)} \sim U[-1, 1]. \quad (7)$$

- Replacement the trajectories (j) in the sample X_0 with probabilities

$$\pi(y^{(j)}, x^{(j)}) = \min \left\{ \frac{p(y^{(j)})}{p(x^{(j)})}, 1 \right\} \quad (8)$$

and obtain the new sample X_1 .

- A sample chain $X_0 \rightarrow X_1 \rightarrow \dots \rightarrow X_m$ has target density distribution function $p(x)$ in the limit.

Markov chain Monte Carlo (Metropolis)

- Disadvantages:
 - High computing costs and time are required;
 - Unable to take into account the physical symmetries.

Alternative approach

- Alternative approach: construction a map $x = g(z)$, $z = g^{-1}(x)$, where z is a set of random variables with a certain density distribution $\pi(z)$.
- Then target distribution $p(\cdot)$ and $g(\cdot)$ are connected as follows:

$$p(x) = \pi(z) \left| \det \frac{\partial z_i}{\partial x_j} \right| = \pi(g^{-1}(x)) \left| \det \frac{dg^{-1}}{dx} \right|. \quad (9)$$

- Neural Network as the map:

$$x^{(j)} = g(z^{(j)}; w), \quad (10)$$

where $z^{(j)} \sim N^n(0, 1)$ and sample $\{x^{(j)}\}$ has target density distribution $p(x)$.

Neural Networks

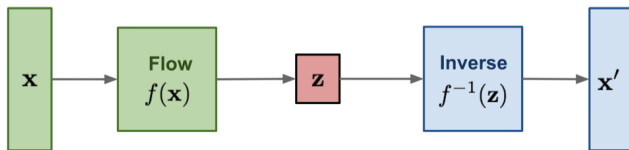


Figure 1 – Normalizing Flow scheme

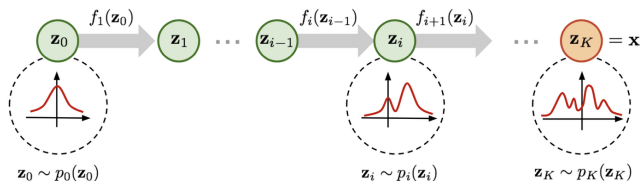


Figure 2 – Normalizing Flow architecture

Normalizing Flows

- The map g is a composition of affine transformations:

$$g = A_n \circ \dots \circ A_1 \quad (11)$$

- We divide z on two parts: $z = u + v$, where, for example, u contains coordinates with even numbers, and v with the odd one.

$$A(u) = u, \quad [A(v)]_k = e^{\theta_{1k}(u)} v_k + \theta_{2k}(u), \quad \theta : \mathbb{R}^{n/2} \rightarrow \mathbb{R}^n. \quad (12)$$

- Loss function reads as follows:

$$\text{Loss}(w) = D_{\text{DL}}(p_g || p) - \ln Z = \int d^n x p_g(x) [\ln p_g(x) + S(x)]. \quad (13)$$

- Orthogonal transform $x = \mathcal{O}g(z)$ was applied to account the shift symmetry of the theory.

Models

- Euclidean action:

$$S_E(x_1, \dots, x_n) = \tau \sum_{i=1}^n [K(x_i - x_{i-1}) + V(x_i)], \quad x_0 = x_n, \quad (14)$$

where $K(\cdot)$ and $V(\cdot)$ are kinetic and potential energies terms.

- Non-relativistic model:

$$H_{\text{kin}} = \frac{p^2}{2m}, \quad K(\xi) = \frac{m\xi^2}{2}. \quad (15)$$

- Relativistic model:

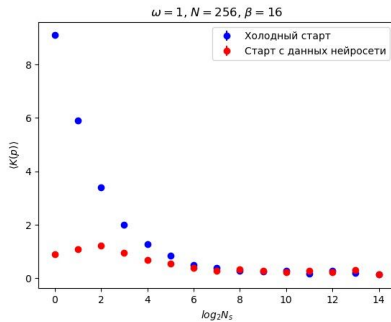
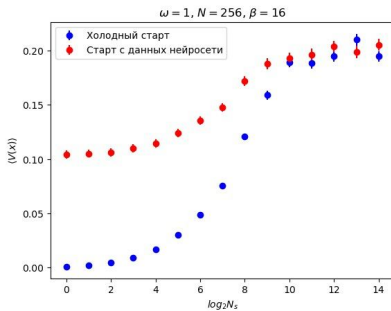
$$H_{\text{kin}} = \sqrt{p^2 + m^2} - m, \quad K(\xi) = \frac{1}{\tau} \ln \left[\frac{m\tau K_1(m\sqrt{\tau^2 + \xi^2})}{\pi\sqrt{\tau^2 + \xi^2}} \right]. \quad (16)$$

- Ultra-relativistic regime $m \rightarrow 0$:

$$H_{\text{kin}} = |p|, \quad K(\xi) = -\frac{1}{\tau} \ln [\pi(\tau^2 + \xi^2)/\tau]. \quad (17)$$

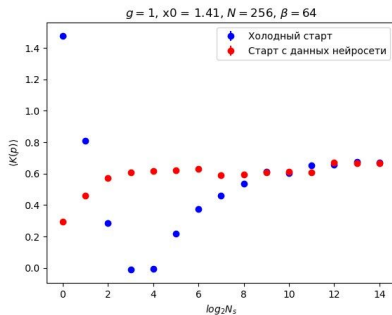
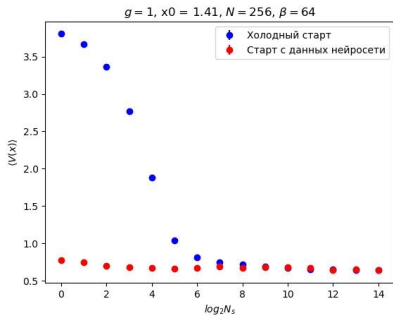
Relativistic Oscillator

$$H = \sqrt{p^2 + m^2} - m + \frac{m\omega^2 x^2}{2} \quad (18)$$



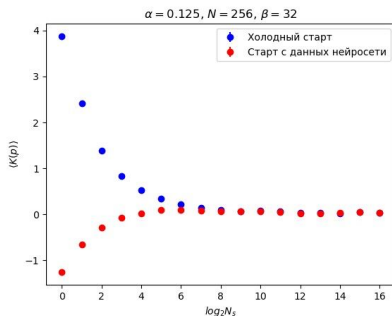
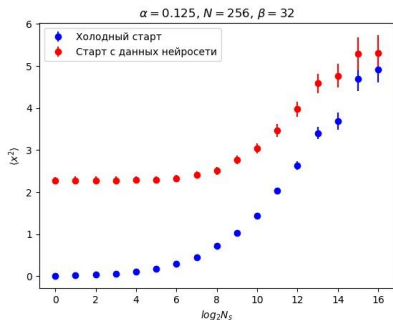
Relativistic double well

$$H = \sqrt{p^2 + m^2} - m + g(x^2 - x_0^2)^2 \quad (19)$$



Relativistic Morse

$$H = \sqrt{p^2 + m^2} - m + \frac{1}{2} \left[(e^{-\alpha x} - 1)^2 - 1 \right] \quad (20)$$



Conclusions

- The use of neural networks makes it possible to speed up the calculation of functional integrals several times.
- The approach is universal: acceleration is observed for different models. This will allow it to be used for a wide range of tasks.
- The symmetry of the problem is taken into account, which may be especially important for applications to the theory of gauge fields.
- The artificial intelligence algorithms used are quite simple. It is expected to significantly improve the results by applying more sophisticated methods.

Thank you!